

P3318 = LAST SECTION

26 APRIL

IN CLASS :  $T = T_{\text{trans}} + T_{\text{rot}}$

$$\begin{aligned} & \uparrow & & \uparrow \\ & \frac{1}{2} M V^2 & & \frac{1}{2} \sum_{\alpha} M_{\alpha} (\vec{\omega} \times \vec{r}_{\alpha})^2 \\ & & & \uparrow \\ & & & = \frac{1}{2} \sum_{ij} I_{ij} \omega_i \omega_j \end{aligned}$$

$$I_{ij} = \frac{\sum_{\alpha} M_{\alpha} \left[ r_{\alpha}^2 \delta_{ij} - (\vec{r}_{\alpha})_i (\vec{r}_{\alpha})_j \right]}{\uparrow} \quad \leftarrow \text{COMPARE TO LEGENDRE DECOMPOSITION}$$

continuous mass distribution

$$\sum_{\alpha} M_{\alpha} \rightarrow \int d^3 \vec{r} \rho(\vec{r})$$

IN PARTICULAR :

$$I = \int d^3 \Omega \rho(\Omega) \begin{pmatrix} y^2 + z^2 & -xy & -xz \\ -xy & x^2 + z^2 & -yz \\ -xz & -yz & x^2 + y^2 \end{pmatrix}$$

can rotate s.t.  $I = \begin{pmatrix} I_1 & & \\ & I_2 & \\ & & I_3 \end{pmatrix}$   
↑ PRINCIPAL AXES                      ↑ PRINCIPAL MOMENTS

[examples from Tong, Marion & Thornton]

eg. the principal moments are  $\mathbb{R}, \geq 0$

Pf. given an arbitrary vector  $\vec{c}$

$$\begin{aligned}\vec{c}^T I \vec{c} &= I_{ab} c^a c^b \\ &= \sum_{\alpha} M_{\alpha} (\vec{r}_{\alpha}^2 \vec{c}^2 - (\vec{r}_{\alpha} \cdot \vec{c})^2) \\ &\geq 0\end{aligned}$$

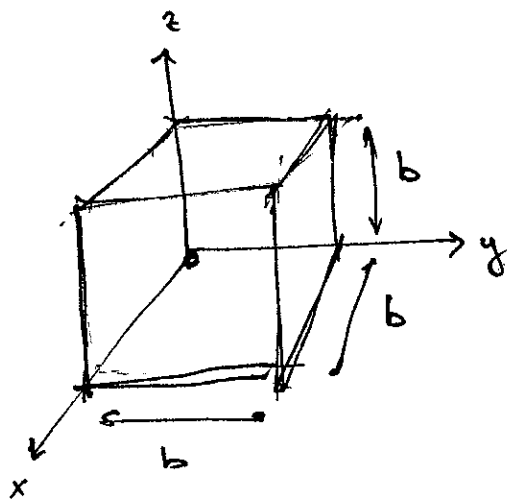
$\vec{c} = 0$  only if all  $\vec{r}_{\alpha}$  ~~collinear~~ <sup>COLLINEAR</sup> w/  $\vec{c}$

SUPPOSE  $\vec{c}$  IS THE  $A^{\text{th}}$  EIGENVECTOR OF  $I$   
THEN:

$$I_{ab} c^a c^b = I_A \vec{c}^2$$

$$\Rightarrow \boxed{I_A \geq 0} \quad \text{---}$$

eg.



HOMOGENEOUS; DENSITY  $\rho$   
TOTAL MASS  $M$

ORIGIN NOT @ CM.  
 $\uparrow$  silly choice, eh?

BEFORE FORCE THE FIRST COUPLE:

$$I_{11} = \rho \int_0^b dz \int_0^b dy (y^2 + z^2) \int_0^b dx$$

$$\underbrace{\int_0^b dy (y^2 + z^2)}_{\frac{1}{3}y^3 + yz^2 \Big|_0^b} \underbrace{\int_0^b dx}_b$$

$$= \rho \int_0^b dz \left( \frac{1}{3}b^3 + bz^2 \right) \cdot b$$

$$= \rho b \cdot \left( \frac{1}{3}b^3 + \frac{1}{3}b^3 \right)$$

$$= \frac{2}{3} \rho b^5 \leftarrow = \frac{2}{3} Mb^2$$

$$I_{12} = -\rho \int_0^b x dx \int_0^b y dy \int_0^b dz$$

$$= -\rho \left( \frac{1}{2}b^2 \right) \cdot \left( \frac{1}{2}b^2 \right) \cdot b$$

$$= \left( -\frac{1}{4} \rho b^5 \right) \leftarrow = -\frac{1}{4} Mb^2$$

phew. should we do the rest?

No: BY SYMMETRY, ALL DIAGONAL ELEM EQUAL  
 & ALL OFF-DIAG ELEM EQUAL.

$$I_{11} = I_{22} = I_{33} = \frac{2}{3} \rho b^5 \quad \checkmark$$

$$I_{12} = I_{13} = I_{23} = -\frac{1}{4} \rho b^5$$

obs: not diagonal... silly coordinates.

## PRINCIPAL AXES

$$\text{WANT: } I_{ij} = I_i \delta_{ij}$$

$$\begin{aligned} \uparrow \text{ then } L_i &= \sum_j I_{ij} \omega_j = I_i \omega_i \\ T_{\text{rot}} &= \frac{1}{2} \sum_{ij} I_{ij} \omega_i \omega_j = \frac{1}{2} \sum_i I_i \omega_i^2 \end{aligned}$$

~~THIS~~ BASIS WHERE ANGULAR VEL & ANGULAR MOMENTUM ARE AMONG THE SAME (PRINCIPAL) AXES.

$$\vec{L} = I \vec{\omega} \quad \leftrightarrow \quad \vec{L} - I \vec{\omega} = 0$$
$$\begin{matrix} \uparrow & \uparrow \\ \begin{pmatrix} I \omega_1 \\ \vdots \end{pmatrix} & \begin{pmatrix} I_{11} \omega_1 + I_{12} \omega_2 + I_{13} \omega_3 \\ \vdots \end{pmatrix} \end{matrix}$$

$$\begin{aligned} (I_{11} - I) \omega_1 + I_{12} \omega_2 + I_{13} \omega_3 &= 0 \\ I_{21} \omega_1 + (I_{22} - I) \omega_2 + I_{23} \omega_3 &= 0 \\ \vdots & \end{aligned}$$

Nontrivial solution:

$$\begin{vmatrix} (I_{11} - I) & I_{12} & I_{13} \\ I_{21} & (I_{22} - I) & I_{23} \\ I_{31} & I_{32} & (I_{33} - I) \end{vmatrix} = 0$$

characteristic polynomial.

moments

Principal ~~axes~~ of cube

$$\begin{pmatrix} \frac{2}{3}\beta - I & -\frac{1}{4}\beta & -\frac{1}{4}\beta \\ -\frac{1}{4}\beta & \frac{2}{3}\beta - I & -\frac{1}{4}\beta \\ -\frac{1}{4}\beta & -\frac{1}{4}\beta & \frac{2}{3}\beta - I \end{pmatrix} = 0$$

$$\begin{pmatrix} (-\frac{11}{12}\beta + 1) & (\frac{11}{12}\beta - 1) & 0 \\ & & \end{pmatrix} = 0$$

$$\begin{pmatrix} -1 & 1 & 0 \\ & & \end{pmatrix} (\frac{11}{12}\beta - 1) = 0$$

Expand :

$$\left( \frac{11}{12}\beta - 1 \right) \left[ \left( \frac{2}{3}\beta - 1 \right)^2 - \frac{1}{16}\beta^2 - \frac{1}{4}\beta \left( \frac{2}{3}\beta - 1 \right) - \frac{1}{16}\beta^2 \right]$$

[work]  $\left( \frac{1}{6}\beta - 1 \right) \left( \frac{11}{12}\beta - 1 \right) \left( \frac{11}{12}\beta - 1 \right) = 0$

PRINCIPAL MOMENTS:

$I_1 = \frac{1}{6}\beta$	} identical
$I_2 = \frac{11}{12}\beta$	
$I_3 = \frac{11}{12}\beta$	

→  $I_1$  is axis of sym

can check plug in  $I = I_1 = \frac{1}{6} B$   
 into system of eq.  $\vec{I} - I\vec{\omega} = 0$

↳ OBTAIN:  $\omega_1^{(1)} = \omega_2^{(1)} = \omega_3^{(1)}$

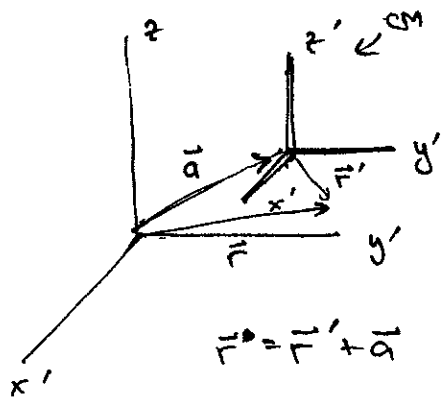
AXIS POINTS IN DIAGONAL.

other two are orb in plane PERP. to this.

~~Rotation to the principal natural basis:  
 transformation of a tensor:~~

~~$I \Rightarrow U I U^T$~~

Parallel axis thm



$$I_{ij} = \sum_{\alpha} m_{\alpha} \left( \delta_{ij} r_{\alpha}^2 - r_{\alpha i} r_{\alpha j} \right)$$

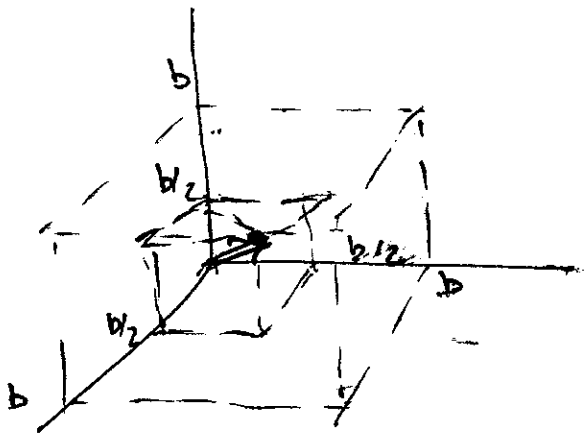
$$\begin{aligned}
 I_{ij} &= \sum_{\alpha} M_{\alpha} \left( \delta_{ij} (\vec{r}'_{\alpha} - \vec{a})^2 + (r'_{\alpha i} + a_i)(r'_{\alpha j} + a_j) \right) \\
 &= I'_{ij} + \left( \sum_{\alpha} M_{\alpha} r_{\alpha}^M \right) \left( \delta_{ij} a^2 - a_i a_j \right) \\
 &\quad + \sum_{\alpha} M_{\alpha} \left( 2 \delta_{ij} \vec{r}' \cdot \vec{a} - a_i r'_{\alpha j} - r'_{\alpha i} a_j \right)
 \end{aligned}$$

$\uparrow$   
 last line vanishes by  $\sum_{\alpha} M_{\alpha} r'_{\alpha i} = 0$   
 since we're in CM frame

$$\boxed{I'_{ij} = I_{ij} - M(a^2 \delta_{ij} - a_i a_j)}$$

$\text{wt cm} \quad \text{cm}$

eg. CUBE w/ coord sys @ origin



PRINCIPAL MOMENT OF INERTIA  
 TENSOR IS  $I'$

$$\begin{cases}
 I'_{11} = I'_{22} = I'_{33} = \frac{5}{3} M b^2 \\
 I'_{12} = I'_{23} = I'_{13} = -\frac{1}{4} M b^2
 \end{cases}$$

$$a_1 = a_2 = a_3 = b/2$$

$$\begin{aligned}
 \text{Then } I_{11} &= I'_{11} - M(a_2^2 - a_1^2) \\
 &= I'_{11} - M(a_2^2 + a_3^2) \\
 &= \frac{2}{3}Mb^2 - 2 \cdot \left(\frac{1}{4}Mb^2\right) \\
 &= \frac{1}{6}Mb^2
 \end{aligned}$$

$$\begin{aligned}
 I_{12} &= I'_{12} - M(-a_1 \cdot a_2) \\
 &= -\frac{1}{4}Mb^2 + \left(\frac{1}{4}Mb^2\right) \\
 &= 0 \quad \checkmark
 \end{aligned}$$

3 others by symmetry.

$$I = \text{diag}(I_{11}, I_{11}, I_{11}) \quad \checkmark$$

$$I \sim \mathbb{1}$$

as long as origin is @ cm,  
no ~~preffer~~ preferred axis

?

$$I \rightarrow UIU^T = I$$



# Dynamics of the falling cat

CONTINUED FROM LAST TIME.

REPRESENTATIVE

$$\vec{L} = \sum M_\alpha \underbrace{\vec{r}_\alpha \times \dot{\vec{r}}_\alpha}_{\vec{L}_\alpha}$$

$$\vec{r}_i = R \tilde{r}_i$$

$$= \sum M_\alpha \left[ (R \tilde{r}_\alpha) \times (R \dot{\tilde{r}}_\alpha) + (R \tilde{r}_\alpha) \times (R \dot{\tilde{r}}_\alpha) \right]$$

falling cat:  $\vec{L} = 0$

[claim:  $R_{ab} = \epsilon_{abc} \tilde{I}_{cd}^{-1} \tilde{L}_d$ ]

$$\tilde{I}_{ab} = \sum_\alpha M_\alpha \left[ \tilde{r}_\alpha^2 \delta_{ab} - (\tilde{r}_\alpha)_a (\tilde{r}_\alpha)_b \right]$$

$$L_a = \epsilon_{abc} \sum_\alpha M_\alpha \left[ R_{bd} R_{ce} (\tilde{r}_\alpha)_d (\dot{\tilde{r}}_\alpha)_e + R_{bd} \dot{R}_{ce} (\tilde{r}_\alpha)_d (\tilde{r}_\alpha)_e \right]$$

mult by  $\epsilon_{afg}$  use  $\epsilon_{abc} \epsilon_{afg} = \begin{pmatrix} \delta_{bf} \delta_{cg} \\ -\delta_{bg} \delta_{cf} \end{pmatrix}$

$$\epsilon_{afg} L_a = \sum_\alpha M_\alpha \left[ R_{fd} R_{ge} (\tilde{r}_\alpha)_d (\dot{\tilde{r}}_\alpha)_e - (\text{DOT MOVE TO } R_{ge}) - R_{gd} R_{fe} (\tilde{r}_\alpha)_d (\dot{\tilde{r}}_\alpha)_e + (\text{--- } R_{fe}) \right]$$

Multiply by  $R_{fb} R_{gc}$

Use:  $R_{fb} R_{fd} = (RR^T)_{bd} = \delta_{bd}$

$$R_{fb} R_{gc} \epsilon_{afg} L_a = \sum_{\alpha} M_{\alpha} \left[ (\tilde{r}_{\alpha})_b (\dot{\tilde{r}}_{\alpha})_c - (\dot{\tilde{r}}_{\alpha})_c (\tilde{r}_{\alpha})_b - R_{bd} (\dot{\tilde{r}}_{\alpha})_c (\tilde{r}_{\alpha})_d + R_{cd} (\tilde{r}_{\alpha})_b (\dot{\tilde{r}}_{\alpha})_d \right]$$

$$= 0$$

~~JUST USE~~

LOOK @ APPARENT ANGULAR MOMENTUM

$$\tilde{L}_a = \epsilon_{abc} \sum_{\alpha} M_{\alpha} (\tilde{r}_{\alpha})_b (\dot{\tilde{r}}_{\alpha})_c$$

$$\downarrow = \cancel{R_{fb} R_{gc} \epsilon_{afg} L_a}$$

$$\tilde{L}_1 = \sum_{\alpha} \left[ (\tilde{r}_{\alpha})_2 (\dot{\tilde{r}}_{\alpha})_3 - (\dot{\tilde{r}}_{\alpha})_3 (\tilde{r}_{\alpha})_2 \right]$$

$$= \sum_{\alpha} M_{\alpha} \left[ R_{21} (\tilde{r}_{\alpha})_3 (\tilde{r}_{\alpha})_1 + R_{23} (\dot{\tilde{r}}_{\alpha})_3 (\tilde{r}_{\alpha})_3 - R_{31} (\tilde{r}_{\alpha})_2 (\tilde{r}_{\alpha})_1 - R_{32} (\dot{\tilde{r}}_{\alpha})_2 (\tilde{r}_{\alpha})_2 \right]$$

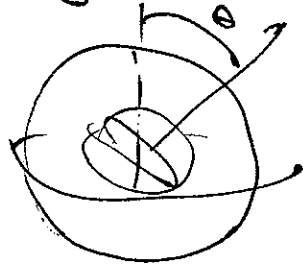
$$= \tilde{I}_{11} R_{23} + \tilde{I}_{12} R_{31} + \tilde{I}_{13} R_{12}$$

$$\tilde{I}_{ab} = \sum_{\alpha} M_{\alpha} \left[ \tilde{r}_{\alpha}^2 \delta_{ab} - (\tilde{r}_{\alpha})_a (\tilde{r}_{\alpha})_b \right]$$

$$= \left[ \frac{1}{2} \epsilon_{abc} \tilde{I}_{1a} R_{bc} \right]$$

A more direct example:

rotating concentric spheres



$\Rightarrow$  CAN CHANGE <sup>REL</sup> ORIENTATION w/o ESTAB ANGIULAR MOMENTUM.

SEE: Gauge kinematics of Deformable Bodies  
Shapere & Wilczek