

## Lecture 7

- calculus of variations applied to mechanics: Hamilton's Principle

Recall that a necessary condition for the time functional  $T$  to be minimized by the function  $\tilde{y}$  is that all the ~~first~~ variational derivatives of  $T$ , at the "point"  $\tilde{y}$ , are zero:

$$\left. \frac{\delta T}{\delta y(x)} \right|_{\tilde{y}} = 0 \quad \text{@ all } x$$

equivalently: 
$$\left. \frac{\partial F}{\partial y} - \frac{d}{dx} \left( \frac{\partial F}{\partial \left( \frac{dy}{dx} \right)} \right) \right|_{y=\tilde{y}} = 0 \quad (1)$$

The structure of this equation is identical to the E-L equations of mechanics. Consider a system with one degree of freedom  $q$ . Now make the identifications,

$$\begin{array}{ccc}
 y(x) & \longleftrightarrow & q(t) \\
 \frac{dy}{dx} & \longleftrightarrow & \dot{q} \\
 \frac{d}{dx} & \longleftrightarrow & \frac{d}{dt} \\
 F & \longleftrightarrow & L
 \end{array}$$

The solution of the E-L equation, let's call it  ~~$y$~~   $q$ , now has another interpretation :

Given a system Lagrangian  $L(q, \dot{q}, t)$ , ~~and~~ and the initial and final values of  $q$  in the time interval  $[0, T]$ , the trajectory (motion) taken by the system,  $\tilde{q}(t)$ , has the property that

$$\delta \int_0^T L(q, \dot{q}, t) dt$$

is at an extremum.

Recall that our condition

$$\left. \frac{\delta S}{\delta q(t)} \right|_{\tilde{q}} = 0 \quad \text{@ all } t$$

does not distinguish between "minimize" or "maximize", only that all the "slopes" 3

are zero. In fact, we will see ~~that~~ that in a system as simple as the 1D harmonic oscillator the extremum is actually a saddle point, having both positive and negative curvature (in the multi-dimensional space).

The principle that a mechanical system takes the trajectory that extremizes the action  $S$  is called Hamilton's principle. Why a mechanical system would "want" something to be extremized is an intriguing question, whose answer had to await the discovery of quantum mechanics.

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