

HOMEWORK 7: Einstein's Equation

COURSE: Physics 208, General Relativity (Winter 2017)

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DUE DATE: **Thursday, March 2** in class. (No class on Tuesday, Feb 28.)

You are required to complete the **Reading Assignment** and **Essential Problems** below. Please let me know if these are too time intensive. No extra problems this week.

Reading Assignment

This week: brush up on the derivation of the Einstein equation from as many perspectives as possible. Start reading about gravitational waves¹.

CAVEAT EMPTOR: as always, do what I mean, not necessarily what I say.

Essential Problems

1 Fluid Mechanics from the Relativistic Ideal Fluid

We introduced the stress energy tensor of the ideal fluid in flat spacetime:

$$T^{\mu\nu}(x) = (\rho(x) + P(x)) U^\mu(x) U^\nu(x) - P(x) \eta^{\mu\nu} , \quad (1.1)$$

where ρ , P , and U are the fluid density, pressure, and 4-velocity, respectively.

The conservation of 4-momentum states that $\partial_\mu T^{\mu\nu} = 0$. From this we can re-derive some of the key equations of fluid mechanics:

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = 0 \quad (1.2)$$

$$\rho [\partial_t + (\mathbf{v} \cdot \nabla)] \mathbf{v} = \nabla P . \quad (1.3)$$

These are the continuity and the Euler equations. The bracketed term in the Euler equation is the so-called convective or material derivative.

Our strategy is to derive these equations in the non-relativistic limit of $\partial_\mu T^{\mu\nu}(x) = 0$. The two equations correspond to projections of this conservation equation onto two subspaces: (1) the one-dimensional subspace along the 4-velocity U^μ and (2) the three-dimensional complement of this space.

- (a) Because U is normalized to unity, the projection operator is simply U^μ itself. Note that since the space is one-dimensional, the projection operator only has one tensorial index. The projection operator onto the three-dimensional complement of U is

$$P^\alpha{}_\mu = \delta^\alpha{}_\mu - U^\alpha U_\mu . \quad (1.4)$$

¹For those that are interested, Cellar Door Books (in Canyon Crest) will be having a book club discussion on *Black Hole Blues*, by Janna Levin, in April. The book describes the LIGO experiment and some of the characters behind it. <http://faculty.ucr.edu/~flipt/physci.html>

Confirm that this is orthogonal to the one-dimensional subspace and that this projection operator is properly normalized ($P^2 = P$).

(b) Prove the following useful result using the fact that U is normalized:

$$U^\mu \partial_\nu U_\mu = 0. \quad (1.5)$$

(c) Show that the projection of $\partial_\mu T^{\mu\nu} = 0$ onto the one-dimensional subspace gives the continuity equation. To do this, calculate $U_\nu \partial_\mu T^{\mu\nu}$. The previous ‘useful result’ may, indeed, be useful. As an intermediate step, you’ll find

$$\partial_\mu (\rho U^\mu) - P(\partial_\mu U^\mu) = 0. \quad (1.6)$$

Take the non-relativistic limit by replacing $U = (1, \mathbf{v})$ with $|\mathbf{v}| \ll 1$ and assume low (non-relativistic) pressure, so that $P \ll \rho$.

(d) Use the projection operator (1.4) to derive the continuity equation from $P^\alpha_\nu \partial_\mu T^{\mu\nu} = 0$. You should find many terms vanishing or canceling. As an intermediate step, you should find:

$$\rho(U \cdot \partial)U^\alpha - \partial^\alpha P + U^\alpha(U \cdot \partial)P = 0. \quad (1.7)$$

Confirm that taking the non-relativistic limit and dropping higher-order terms in $\mathbf{v}$ and P gives the Euler equation, above.

Observe that the Euler equation is the generalization of $\mathbf{F} = m\mathbf{a}$ for a fluid, where the “acceleration” is defined as a convective derivative of the 3-velocity $\mathbf{v}$. The simple interpretation for this is that the [total] derivative of the velocity should be expanded:

$$\frac{d\mathbf{v}}{dt} = \frac{\partial \mathbf{v}}{\partial t} + \frac{dx^i}{dt} \frac{\partial \mathbf{v}}{\partial x^i}, \quad (1.8)$$

where the second term is simply $(\mathbf{v} \cdot \nabla)\mathbf{v}$.

This expression comes up often in the theory of structure formation and galaxy evolution. See, e.g. [astro-ph/9410043](https://arxiv.org/abs/astro-ph/9410043) or your favorite more modern reference². In this context, the “ideal fluid” that is being described by $T^{\mu\nu}$ is dark matter.

2 Covariant Constancy of the Einstein Tensor

In lecture 12 we “hacked” together the Einstein equation. Our strategy was to find something which satisfies the heuristic form (inspired by electrodynamics):

$$(\text{curvature}) = (\text{coupling}) (\text{source}). \quad (2.1)$$

We identified the source as $T_{\mu\nu}$ and the coupling as something proportional to the Newton constant G . We then use the symmetry and two-indexed-ness of $T_{\mu\nu}$ to constrain what the left-hand side might be. After a little thought, we only came up with two possible terms so that the left-hand side must be a linear combination:

$$(\text{prefactor}) (R_{\mu\nu} + \alpha R g_{\mu\nu}). \quad (2.2)$$

²I like <http://www.damtp.cam.ac.uk/user/db275/Cosmology/Chapter4.pdf>.

The prefactor may be absorbed into how we define the coupling on the right-hand side. We fixed the relative size α by requiring that the left-hand side is covariantly constant, $D_\mu(R_{\mu\nu} + \alpha R g_{\mu\nu}) = 0$. This is required since the right-hand side is covariantly constant by the conservation of 4-momentum: $D_\mu T^{\mu\nu} = 0$. (This is, of course, the curved space generalization of the main equation used in problem 1.)

In lecture, we claimed that covariant constancy implied $\alpha = -1/2$ and defined the Einstein tensor as

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}. \quad (2.3)$$

Prove that $D_\mu G^{\mu\nu} = 0$.

HINT: As an intermediate step, you should prove the Bianchi identity:

$$D_\lambda R_{\mu\nu\alpha\beta} + D_\nu R_{\lambda\mu\alpha\beta} + D_\mu R_{\nu\lambda\alpha\beta} = 0. \quad (2.4)$$

How should you do this? Well, it turns out you might want to start by proving the Jacobi identity, which you may remember from quantum mechanics or group theory:

$$[D_\lambda, [D_\mu, D_\nu]] + [D_\nu, [D_\lambda, D_\mu]] + [D_\mu, [D_\nu, D_\lambda]] = 0. \quad (2.5)$$

Then use a result that we showed when we learned about the Riemann tensor:

$$[D_\alpha, D_\beta] A^\mu = R^\mu_{\lambda\alpha\beta} A^\lambda. \quad (2.6)$$

You can, in turn, prove this using $D_\alpha D_\beta A^\mu = D_\alpha(\partial_\beta A^\mu + \Gamma^\mu_{\beta\lambda} A^\lambda)$ and the definition of the Riemann tensor, $R^\mu_{\lambda\alpha\beta} = \partial_\alpha \Gamma^\mu_{\lambda\beta} - \partial_\beta \Gamma^\mu_{\lambda\alpha} + \Gamma^\mu_{\nu\alpha} \Gamma^\nu_{\lambda\beta} - \Gamma^\mu_{\nu\beta} \Gamma^\nu_{\lambda\alpha}$ ($\alpha \leftrightarrow \beta$).

3 Einstein Equation for Schwarzschild

Confirm that Einstein's equation is satisfied for the Schwarzschild metric. HINT: start by showing that it is sufficient to show that $R_{\mu\nu} = 0$. What should you use for $T_{\mu\nu}$? Feel free to use a *Mathematica* package to calculate any curvature quantities.

4 Einstein Equation for Dark Energy

Suppose you have a universe with a spherically symmetric, but time-dependent metric:

$$ds^2 = dt^2 - a(t)^2 d\mathbf{x}^2. \quad (4.1)$$

Fill the universe with a cosmological constant so that the 'matter' action is:

$$S_{\text{stuff}} = - \int d^4x \sqrt{g} \Lambda. \quad (4.2)$$

where $g = -\det g_{\mu\nu}$. Use Einstein's equation to determine the form of the **scale factor**, $a(t)$ up to initial conditions. Feel free to use a *Mathematica* package to calculate any curvature quantities.