

HOMEWORK 2: Adding Velocities & Equivalence Principle

COURSE: Physics 208, General Relativity (Winter 2017)

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DUE DATE: Tuesday, January 31 in class... or, you know, like... whenever.

You are required to complete the **Reading Assignment** and **Essential Problems** below. Please let me know if these are too time intensive. You are invited to explore the ‘extra’ problems as they apply to your goals for this course: **Mathematical Problems** develop geometric intuition, while **Phenomenological Problems** are applications of relativity.

This week’s sound track: “Free Falling” by Tom Petty¹. As you now understand, gravity is simply a consequence of the free fall of an inertial frame in a curved spacetime.

Reading Assignment

Read the following topics. You may choose to read the analogous topics in an appropriate textbook or reference of your preference. Most of this reading is meant to be complementary to the approach in the lectures. For those who would like a solid reference for the material in the lectures, a good place is Weinberg (*Gravitation and Cosmology*, not the newer *Cosmology* book), chapter 3 and the beginning of 4.

- Read chapters 6, 7.1–7.5 of Hartle (at the level of detail that interests you). This connects much of what we’ve been discussing to actual observations.
- Read chapter 8 of Hartle on geodesics, we take a complementary approach in class.

Essential Problems

1 Velocity Addition and Causality

Despite the paradoxes of special relativity, one fundamental tenet of physics is **causality**: if event A can affect event B , then there is *no* reference frame in which B occurs before A . One way to see that special relativity does not violate causality is to check the famous velocity addition formula to confirm that velocities can never add to be greater than the speed of light, $c = 1$.

By “famous” we mean that this addition formula is famously a pain to derive. By “addition” we mean the scenario where we are standing at a train platform watching a **red train** pass by with constant velocity. On this train, someone throws a **blue racquetball** forward so that it, too, has a constant velocity along the train’s motion. We seek to relate the measured racquetball velocity in the frame of the red train to that measured from the platform.

While the crux of gravitation is *differential* geometry, we will solve this problem using high school plane geometry. We follow the derivation from Sander Bais’ excellent popular science book, *Very Special Relativity: An Illustrated Guide*.

¹https://youtu.be/T3phscjgc_A

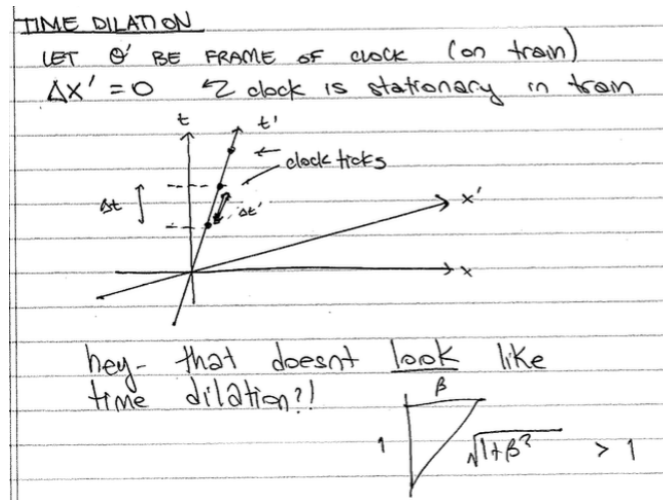
$$\text{XVII} \quad \begin{array}{ccccccc} & + & : & +1 & - & 1 & : \\ & & & & & & c+1 \\ & & & & & & 1+c \\ & & & & & & : \\ & & & & & & c \\ & & & & & & c+1 \\ & & & & & & 1+c \\ & & & & & & : \\ & & & & & & 2H \\ & & & & & & +1 \\ & & & & & & : \\ & & & & & & 1 \end{array}$$

- (j) Check that the diagram confirms this result for $v = 1/2$ and $u' = 1/2$, $u = 4/5$.

- (k) What happens as $u' \rightarrow 1$? Argue that superluminal velocities cannot be generated by velocity addition. This means that physical processes can never cross 'light-cone,' in any frame, and hence that causality is preserved.

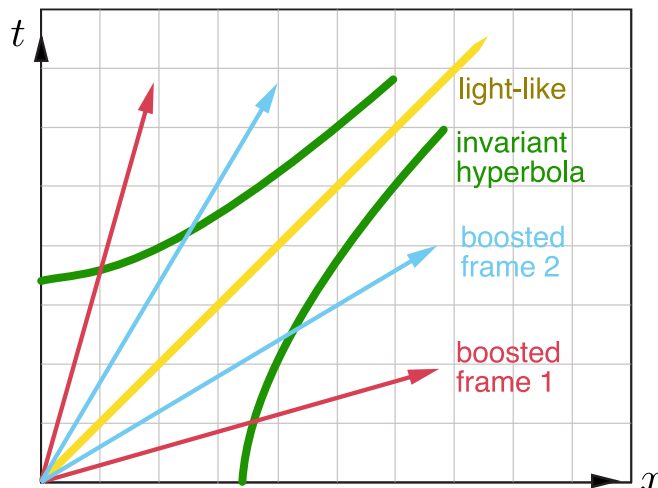
2 Invariant Hyperbolae

Recall in our first lecture that there was an apparent paradox: if space and time are being treated the same, why was it that time is *dilated* while length is *contracted* in a boost? We posted the problem geometrically as follows (from the lecture 1 notes):



Based on simple geometry, it seems like ticks on a light clock are longer in the boosted (primed) frame versus the stationary frame (unprimed). In other words, it looks like time should be contracted. We argued that actually, the Lorentz transformation stretches these ticks. There's an easier way to see this using invariant hyperbolae. In \mathbb{R}^2 , constant radial length corresponds to an invariant circle: $x^2 + y^2 = r^2$. No matter how one rotates the axes, the radius of the circle is preserved. By comparison, in 2D Minkowski space, the invariant interval is given by a hyperbola, $t^2 - x^2 = s^2$.

Using the following figure (adapted again from Bais) and this notion of invariance, argue that time is indeed dilated rather than contracted.



3 Christoffel Symbols of the Sphere (Hartle 8-2)

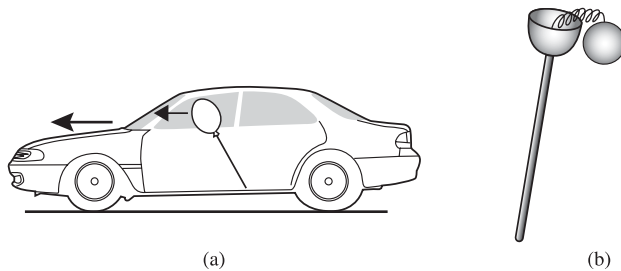
The metric on the sphere of radius a in spherical coordinates is

$$ds^2 = a^2 (d\theta^2 + \sin^2 \theta d\phi^2) . \quad (3.1)$$

- (a) Calculate the Christoffel symbols for this space.
- (b) Show that the great circle (the equator) is a solution to the geodesic equation. **HINT:** Use the freedom to orient the coordinates so that the equation of a great circle is simple.

4 Equivalence Principle Thought Experiments

This is from Ta-Pei Cheng's *Relativity, Gravitation, and Cosmology: A Basic Introduction*, problem 4.2 (including the image below). Here are two 'brain teasers' for the equivalence principle.



- (a) **Forward leaning balloon.** Use the equivalence principle to explain the observation that a helium balloon leans *forward* in a forward-moving car.
- (b) **Old Man's Toy².** On Einstein's last birthday, Eric Rogers gave him a toy composed of bowl with a spring connecting it to a ball. The entire contraption is on a stick. It is actually rather tricky to try to put the ball into the bowl by holding only the bottom of the stick due to the Hooke force from the spring. Use the equivalence principle to suggest an efficient strategy to put the ball into the bowl without directly handling the ball or the bowl.

Phenomenological Problems

5 Particle Accelerator (Hartle 5-12)

The Stanford linear accelerator (part of what is now called the SLAC National Accelerator Laboratory) is an electron–positron collider that is an older cousin of the Large Hadron Collider. In its heyday as a particle physics center, it accelerated electrons from rest to 40 GeV over 2 miles. Idealize the acceleration mechanism as a constant electric field, \mathbf{E} , along the accelerator line and assume an equation of motion

$$\frac{d\mathbf{p}}{dt} = e\mathbf{E} , \quad (5.1)$$

²This is also a title of a popular book on physics by Tony Zee.

for three-momentum \mathbf{p} and electric charge e .

- (a) Assuming the electron starts from rest, find its position along the accelerator as a function of time in terms of its rest mass m and $F = e|\mathbf{E}|$.
- (b) What value of $|\mathbf{E}|$ is necessary to accelerate the particle to its final energy of $E = 40$ GeV.

Use the three-velocity, $\mathbf{v} = \mathbf{p}/E$ and the relativistic version of the famous Einstein relation,

$$E^2 = m^2 + \mathbf{p}^2 . \quad (5.2)$$

In funny units, the answer is $|\mathbf{E}| = 1.2 \times 10^7$ volts per meter. Use:

$$e = 1.6 \times 10^{-19} \text{ C} \quad \text{GeV} = 1.6 \times 10^{-10} \text{ J} \quad \text{meter} = 1610 \text{ mile} . \quad (5.3)$$

6 The Equivalence Principle (Hartle 6-7)

Consider the following change of coordinates from ‘usual’ Cartesian coordinates (unprimed) to funny coordinates (primed),

$$t = (g^{-1} + x') \sinh(gt') \quad x = (g^{-1} + x') \cosh(gt') - g^{-1} \quad y = y' \quad z = z' , \quad (6.1)$$

where g has dimensions of acceleration.

- (a) Transform the line element, ds^2 of ordinary special relativity into the line element of the primed coordinates. ANSWER:

$$ds^2 = (1 + gx')^2 dt'^2 - dx'^2 - dy'^2 - dz'^2 . \quad (6.2)$$

- (b) Assume $gt' \ll 1$. By Taylor expanding $t(t', x')$ and $x(t', x')$, show that the funny variables are simply a uniformly accelerated frame in Newtonian mechanics. Observe how this ‘looks like’ gravity.
- (c) Show that a clock at rest in the primed frame at position $x' = h$ runs faster than a clock at rest at $x' = 0$ by a factor of $(1 + gh)$. Observe that this result in an accelerated frame (but no ‘gravity’) is precisely the same thing we observed in class when we considered a clocks in a gravitational field.

7 Rotating Frames (Hartle 8-4)

The line element of a flat spacetime in a frame that is rotating with angular velocity Ω about the z -axis of an inertial frame is

$$ds^2 = [1 - \Omega^2(x^2 + y^2)] dt^2 - 2\Omega(y dx - x dy)dt - dx^2 - dy^2 - dz^2 . \quad (7.1)$$

- (a) Verify that this matches the ordinary Minkowski metric in spherical coordinates,

$$ds^2 = dt^2 - dr^2 - r^2 d\phi^2 - r^2 \sin^2 \theta d\phi^2 , \quad (7.2)$$

with the substitution $\phi \rightarrow \phi - \Omega t$.

- (b) Find the **geodesic equations** for x , y , and z in the rotating frame.
- (c) Show that in the non-relativistic limit, these reduce to the usual equations of Newtonian mechanics for a free particle in a rotating frame exhibiting the centrifugal and Coriolis force. Recall that, for example, the x -component of the centrifugal force (with $\mathbf{\Omega} = \Omega \mathbf{e}_z$) is $\mathbf{\Omega} \times (\mathbf{\Omega} \times \mathbf{x})$, and that the x -component of the Coriolis force is $2\mathbf{\Omega} \times (d\mathbf{x}/dt)$.

REMARK: If I have sign errors, please fix them. (Do what I *mean*, not what I say.) I somewhat regret asking this question since it's computationally tedious.

Mathematical Problems

8 Poincaré Half Plane

This question is shamelessly borrowed³ from a discussion Zee's *Einstein Gravity in a Nutshell*, chapter I.5. The Poincaré half plane is a funny manifold. It is composed of points (x, y) restricted to $y > 0$ and has the funny metric:

$$ds^2 = \frac{dx^2 + dy^2}{y^2} . \quad (8.1)$$

- (a) Draw the Poincaré half plane as the Cartesian plane with no negative y axis. Now draw an apple somewhere on the plane. Draw the same apple at a different y value, taking rough account for the rescaling of the apparent size coming from the funny metric.
- (b) Does the Poincaré half plane have a boundary? For our purposes, a metric space 'has a boundary' if the "boundary" is a finite distance away from any given point.
- (c) What is the length of the 'straight line' path from $(0, y_*)$ to (x_*, y_*) , for $x_*, y_* \neq 0$? This straight line path is the one given by integrating ds only in the dx direction.
- (d) A **geodesic** is a path of minimal length. Write $ds^2 = L^2 dy^2$, where

$$L = \sqrt{\frac{1 + \left(\frac{dx}{dy}\right)^2}{y^2}} . \quad (8.2)$$

Then the length of a path $x(y)$ is given by $\int ds = \int L dy$. You can solve this as a variational problem, treating y as a 'time' variable and L as a 'Lagrangian':

$$\frac{d}{dy} \frac{\delta L}{\delta(dx/dy)} = \frac{\delta L}{\delta x} . \quad (8.3)$$

Since $\delta L/\delta x = 0$, $\delta L/\delta(dx/dy)$ is an integral of motion. Solve $\delta L/\delta(dx/dy) = 1/b$ for $x(y)$ to find that

$$x - x_0 = \pm \sqrt{b^2 - y^2} , \quad (8.4)$$

so that the geodesics are actually a very simple shape. Draw a geodesic on the Poincaré half plane, label x_0 and b .

³It turns out to also be Hartle, problem 8-12.